Today:

## Recap on 1<sup>st</sup> lecture:

- Basic properties of the ISM
  - HII regions, dark nebulae, reflection nebulae
  - Cold, warm and hot phases of the ISM (properties etc.)
- Look in detail at HII regions:
  - Ionization/recombination, emission/absorption
  - Thermal balance (heating/cooling)
  - Sizes

#### Vela Supernova Remnant





Data from the Digitized Sky Survey Image processing by Davide De Martin

#### ISM:

built from remains of stars
(i) during life: stellar winds (slow and steady during lifetime)
(ii) end of life: elected massive envelopes (PN\_novae)



## Can we classify the ISM?

 <u>Diffuse emission nebulae</u>: When in the vicinity of OB-type stellar associations (i.e. near young stars), they become <u>ionized</u> are called <u>HII</u> <u>regions</u>:

- The gas emits photons due to recombination.
- Generally found near the plane of the Galaxy, at heights of  $\geq$  50 pc.
- Masses ~ 100–10000 M<sub>☉</sub>; Sizes ~ few pc; Temps ~ 10000 K; Densities ~ 10<sup>3</sup> hydrogen ions/cm<sup>3</sup>
   (compare with air density of 10<sup>19</sup>/cm<sup>3</sup>, or stellar atmosphere ~ 10<sup>15</sup>/cm<sup>3</sup>)

• <u>Dark clouds</u> are denser and colder  $>10^4$ /cm<sup>3</sup> and T ~ 10–100K - are potential sites of <u>star-formation</u>. They block most optical radiation.

 $H_2$  molecules form within them on the surface of dust grains. Variety of sizes: pc-sized up to <u>Giant Molecular Clouds</u>

 <u>Reflection nebulae</u> near bright stars appear bluish due to the efficient scattering, by dust particles, of blue-wavelength light (say 400 nm) – same process as that which makes the sky blue

#### Distribution of Cold Gas in the ISM





Distance = 450 pc Diameter ~ 1° (similar to the full moon)

Composition:

- Ionized gas:
- H ~ 90% (by number)
- He ~ 10% (by number)
- <1% Heavier elements

• Dust

### The Trapezium Cluster



dominant star:  $\theta^1$  Ori-C spectral type O6V: T ~ 20 000 K

emits a great number of energetic photons that ionize the gas



Bok globules

relatively much denser that general diffuse ISM clouds:

densities ~10<sup>4</sup> - 10<sup>9</sup> particles per cm<sup>3</sup>

- made up of atoms, molecules and <u>dust particles</u>
- have low temperatures of 10-100 K.



Gamma CrA

Dark Nebula Bernes 157

NGC 6729

**Double Star BRS 014** 

**Reflection Nebula** IC 4812

NGC 6723

**Reflection Nebula** NGC 6726/6727

Variable Nebula

Epsilon CrA



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#### ISM: average properties

 Mean density of ISM (in our Galaxy) ~ 10<sup>6</sup> particles/m<sup>3</sup> (1 particle/cm<sup>3</sup>) but is HIGHLY CLUMPED: filling factor ~3%.

• The GAS is mainly:

Hydrogen (90%) Helium (10%) metals (C,N,O, ... Fe) <1% by number

- Dust grains : gas atoms ratio = approx 1 : 10<sup>12</sup>
- Dust grains have a power-law size distribution radii ~  $5 \times 10^{-9} - 2 \times 10^{-7}$  m (smaller grains are more abundant)

## A global model: the 3+1 phases of the ISM

• Cold, neutral medium (CNM)

 $n \sim 1-10^3$  /cm<sup>3</sup>, T < 100 K, volume fraction:  $\sim 1-5\%$ 

- Hot, ionized medium (HIM)
   *n* ~ 10<sup>-4</sup>-10<sup>-2</sup> /cm<sup>3</sup>, *T* ~ 10<sup>6</sup>-10<sup>7</sup> K, volume fraction: ~30-70%
- Warm interface media
  - Warm ionized medium (WIM)

 $n \sim 0.01 \text{ /cm}^3$ ,  $T \sim 1000 \text{ K}$ , volume fraction:  $\sim 20-50\%$ 

hydrogen ionization fraction ( $X_{\rm H}$ ) ~ 70%

 Warm neutral medium (WNM) n ~ 0.1–10 /cm<sup>3</sup>, T ~ 1000–5000 K, volume fraction: ~10–20% hydrogen ionization fraction (X<sub>H</sub>) ~ 10%

Dark neutral (molecular) clouds

 $n \sim 10^3 \Box 10^6$  /cm<sup>3</sup>,  $T \sim 10 \Box 50$  K, volume fraction: <1%

## A global view





FIG. 2

### Hot, ionized medium (HIM)

Arises through the heating and energy input from overlapping Supernova explosions and their remnants: directly seen in X-rays.



Optical = red and green X-ray = blue



### Hot, ionized medium (HIM)

Arises through the heating and energy input from overlapping Supernovae explosions and their remnants: directly seen in X-rays.



- Supernovae occur in the Galactic plane: causes hot gas to rise to high distances above/below the plane of the Galaxy
- Gives rise to a halo of hot gas around the galaxy.
- This gas then cools and falls back to the galactic plane, and is replenished by further SN: the Galactic Fountain model.

### Cold, neutral medium (CNM)

• dominated by diffuse clouds with  $n \sim 10^{1}-10^{3}$  /cm<sup>3</sup>,  $T \sim 30-100$  K, individual cloud radii = few pc

• cold enough that simple molecules can form, e.g.  $H_2$  and CO.

• <u>however</u>, most INTERSTELLAR MOLECULES are mainly found in the denser DARK CLOUDS (Barnard Objects and Bok Globules) with  $n \sim 10^6-10^7$  /cm<sup>3</sup> and  $T \sim 10$  K.

• Still denser regions are the GIANT MOLECULAR CLOUDS with  $n \sim 10^{10}$  /cm<sup>3</sup> and  $T \sim 50$  K. These are associated with sites of new STAR FORMATION.



# THE MILKY WAY PROJECT

### www.milkywayproject.org



## Today:

## HII regions:

- Ionization/recombination, emission/absorption
- Thermal balance (heating/cooling)
- Sizes
- The addition of metals
  - Forbidden lines
  - Metal line cooling
  - Density/temperature diagnostics



Rosette Nebula

### 2 O-stars



## Photoionisation

Photoionisation: the removal of electrons from gas atoms by photons.

H + (ionizing photon)  $\rightarrow$  p + e<sup>-</sup> (Eqn. 1)

## Photoionisation



13.6 eV = λ < 912 Å **Photoionisation**: the removal of electrons from gas atoms by photons.

H + (ionizing photon)  $\rightarrow$  p + e<sup>-</sup> (Eqn. 1)

For Hydrogen (most abundant atom) photons need to have an energy of *at least* **13.6 eV** (corresponding to  $\lambda < 912$  Å)

13.6 eV therefore = the "<u>lonisation Potential</u>" of Hydrogen (*I*<sub>H</sub>) = called **one Rydberg** 

Only stars with T > 20000 K (spectral types O and B) emit 'ionising' photons

## Recombination

Attraction between protons and electrons leads to the recapture of the electron.

 $p + e^- \rightarrow H + (ionizing photon)$ 

### **Recombination cascade in H-atom**







Planetary Nebula: NGC 3242 (ESO 1.5-m in Chile)

Blue = recombination lines of H and He Red = forbidden lines of 'metals'

To a good approx. the nebula contains only neutral H, protons (p), and electrons (e<sup>-</sup>)

The energy of the free electrons can be "anything" above 13.6 eV

The *total* energy available for photoionization depends on the *total* number of photons.

This is a function of the effective temperature ( $T_{eff}$ ) of the ionizing star

The energy is spread out (**shared**) between all the particles (by collisions) and acquires a single '<u>kinetic temperature</u>' for *all 3 types* of particle. We say the *plasma* becomes '<u>thermalized</u>'.

The distribution of velocities within the gas can now be described by a <u>'Maxwellian distribution</u>'

## The Maxwellian Distribution



Photoionisation rate [number of events/m<sup>3</sup>/sec]:

 $N_{\rm PI} = a n_{\rm HI} J$  [m<sup>-3</sup> s<sup>-1</sup>]

**a** = photoionsation cross section (=  $6.8 \times 10^{-22} \text{ m}^2$ ) **n**<sub>HI</sub> = number density of neutral H (HI) **J** = rate of incident ionising photons (with E >13.6 eV) (m<sup>-2</sup> s<sup>-1</sup>)



#### cross section = likelihood (probability) of interaction between particles

Photoionisation rate [number of events/m<sup>3</sup>/sec]:

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#### Recombination rate [number of events/m<sup>3</sup>/sec]:

$$N_{\rm R} = n_{\rm e} n_{\rm p} \alpha_n = n_{\rm e}^2 \alpha_{\rm R}$$
 [m<sup>-3</sup> s<sup>-1</sup>]

 $n_{\rm e}$ ,  $n_{\rm p}$  = number density of electrons & protons (in equilibrium  $n_{\rm e} = n_{\rm p}$ )  $\alpha_{\rm R}$  = recombination coefficient = 2×10<sup>-16</sup>  $T_{\rm e}^{-3/4}$  m<sup>3</sup> s<sup>-1</sup>]
# **Ionisation** balance

Equilibrium occurs when there is a balance between the forward and backward rates

i.e. Photoionization rate,  $N_{PI}$  = Recombination rate,  $N_{R}$ 

H + (ionising photon)  $\leftrightarrow$  p + e<sup>-</sup>

The degree of ionisation can be described as:

 $n_{\rm e} = n_{\rm p} = X n$ 

*X* = fraction of *ionised* H atoms
 *n* = number density of H nuclei (proton+neutral atom densities)

 $n_{\rm HI} = n - X n$ 

 $n_{\rm HI}$  = number density of neutral H (HI)

Now we can use these equations in the previous relations...

- $N_{PI} = a n_{HI} J$ •  $N_{R} = n_{e}^{2} \alpha_{R}(T_{e})$
- $n_{\rm e} = n_{\rm p} = X n$
- $n_{\rm HI} = n X n$

In equilibrium,  $N_{PI} = N_R$ , so:  $a (n - X n) J = (X n)^2 \alpha_R$ 

We can now calculate X (the ionisation fraction) for a specific value of J and n at some distance from a star, r

$$J = S_* / (4 \pi r^2) \qquad [m^{-2} s^{-1}]$$

 $S_*$  = no. of ionising photons emitted by the star per sec

For a typical O-star (e.g. the O6V star in the Trapezium)  $S_* = 10^{49} \text{ s}^{-1}$ 

Typical nebula,  $n = 10^8 \text{ m}^{-3}$ ;  $T_e = 10^4 \text{ K}$ 

At r = 1 pc,  $J = 8.4 \times 10^{14}$  m<sup>2</sup> s<sup>-1</sup>

giving X = 0.9999999i.e. the HII region is almost fully ionised!





FIGURE 2.4 Ionization structure of two homogeneous  $H + He \mod H$  II regions.





- 1. no. of ionising photons
  - Hotter star
  - More hot stars
- 2. amount of gas (gas density)

IC 1274





## Thermal balance in ionized nebulae

Photoionization adds kinetic energy (KE) into the H cloud since it creates hot, fast photoelectrons.

*T* can increase indefinitely though... Energy must be lost somehow

Energy can be lost through photons: H recombination means KE of  $e^- \rightarrow$  photon which can escape

On average, heat lost per recombination =  $3/2 k T_e$  so Cooling rate is:

 $L = (3/2 \ k \ T_e) N_R$  [J m<sup>-3</sup> s<sup>-1</sup>]

Energy input, or the Heating rate

 $G = N_{PI} Q$  [J m<sup>-3</sup> s<sup>-1</sup>]

**Q** = the heating energy injected into the gas per photoionization (in Joules)

in equilibrium, G = L (and  $N_R = N_{Pl}$ )  $\rightarrow$   $T_e = (2/3)Q/k$ .

However, Q per ionization is: (3/2)  $kT_* \rightarrow T_e = T_*$ For OB-type ionizing stars  $T_* = 30\ 000 - 60\ 000$  K These temperatures are much higher than observations Typical nebula gas temperature has  $T_e \sim 10\ 000$  K.

#### → There must be other ways to cool the gas

What if we relax our assumption of a pure H nebula?

# Metals

- Real nebulae contain more than just H
- Metals (heavy elements, C, N, O, Ne, S, Ar, Cl, etc.) are found in proportion to Hydrogen of  $\sim 10^{-4} 10^{-8}$





H recombination = ALLOWED transitions

high probability of occurring (~10<sup>9</sup> per sec)

#### excitation/de-excitation of bound electrons





energy levels (→ transitions) in metals are not as simple as H!

$$n = 3$$

$$n = 2$$

$$n = 1$$

*n* = 1

#### Collisionally excited lines (CELs)



# **Forbidden lines**

Forbidden line: arises when an electron is excited by a collision into a 'metastable state'.

When densities are higher (e.g. our atmos. – greater than about  $10^8$  per cm<sup>3</sup>), electron would almost immediately be knocked out of metastable state by collision and not be given time to emit a photon (collisional de-excitation = no photon).

But in a low density nebula, the time between collisions ~ 10 - 10,000 seconds

- (v. long time)
- → allowed time to radiate spontaneously

In a nebular environment, practically every ion goes to the ground state by forbidden radiation

The term 'Forbidden' refers to quantum mechanical rules - a little misleading!

More intuative name: <u>Collisionally excited lines</u> (CELs) - line emission following collision with electron



CELs (forbidden lines) are written with [] e.g. [OII] 3726Å



critical density for collisional de-excitation =  $7 \times 10^5$  cm<sup>-3</sup>

 $10^5$  cm<sup>-3</sup> = density where collisional deexcitation of many of the bright lines begins to matter





## Forbidden line cooling

transition rates very low:  $\sim 10^{-3} - 100$  per sec (compared to H recombination rates:  $10^9$  per sec)

Therefore photons are v. likely to escape nebula before being absorbed

 $\rightarrow$  can ignore absorption

They can remove a lot of heat from the nebula, thus solving our problem

Common forbidden lines: <u>Optical</u>: [OIII] 4959,5007 Å, [NII] 6548,6584 Å, [SII] 6717,6731 Å <u>Infrared</u>: [OIII] 52,88 μm, [NIII] 57 μm



Consequence of forbidden-line cooling:

the higher the metallicity (i.e. heavy-element content) of a nebula, the faster it cools to thermal equilibrium

and the stronger the forbidden lines are



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Topical podcast with two monthly episodes, featuring interviews, news items, night sky, "Ask an astronomer", humour (and occasionally videos)

## Temperatures



Observation  $\rightarrow$  T<sub>e</sub> of ~7000–10000 K for HII regions ~9000–15000 K for PNe



### **Densities**



Observations  $\rightarrow n_e \sim 100 \text{ cm}^{-3}$  for HII regions  $\sim$  a few 1000 cm<sup>-3</sup> for PNe



Intensity ratio

Electron density (cm<sup>-3</sup>)

## Abundances

Now we know  $T_{\rm e}$  and  $n_{\rm e}$ 

we can calculate the <u>abundance</u> of all the different ions emitting the forbidden lines

For example in the case of oxygen ions (but also for ions of C, N, Ne, S, Ar, Cl etc):  $\frac{N(O^{2^+})}{N(H^+)} \sim \frac{n_e}{f(T_e)} \frac{I(/4959)}{I(Hb)}$ 

 $n_e$  = electron density  $f(T_e)$  = fraction of O<sup>2+</sup> ions able to emit at 4959 Å (strong dependence on nebular temp)  $I(\lambda 4959)/I(H\beta)$  = flux of the [OIII] 4959 Å line relative to H $\beta$ 

→ Gives the number of O<sup>++</sup> ions *relative to* the number of H<sup>+</sup> (the most abundant ion in a nebula).



Measure forbidden lines from <u>all</u> ionic stages of an element (e.g. O, O<sup>+</sup>,  $O^{2+}$ )

add up all the abundances to find the <u>total abundance</u> relative to hydrogen.

#### Problem:

We need to know abundance of <u>all</u> ions of particular element to really know its abundance - not always possible

some lines are weak→ need to use big telescope

For more distant objects (e.g. galaxies), this becomes even more of a problem!

some lines are unobservable

→ ionization corrections (theoretical)

## End of last lecture

# Last lecture(s):

Recap on essentials from last lecture Looked in detail at: HII regions

- Photoionization and recombination
- Heating and cooling processes
  - H heating problem → Metals → Forbidden lines
  - Temp and density diagnostics, abundances









HII regions: Practical methods & examples

#### Cooler phases of ISM - non-ionised material





## Temperatures



Observation  $\rightarrow$  T<sub>e</sub> of ~7000–10000 K for HII regions ~9000–15000 K for PNe

### **Densities**



Observations  $\rightarrow n_e \sim 100 \text{ cm}^{-3}$  for HII regions  $\sim$  a few 1000 cm<sup>-3</sup> for PNe

Now we know  $T_{\rm e}$  and  $n_{\rm e}$ 

we can calculate the <u>abundance</u> of all the different ions emitting the forbidden lines

add up all the abundances to find the <u>total abundance</u> relative to Hydrogen.

#### Problem:

We need to know abundance of <u>all</u> ions of particular element to really know its abundance - not always possible

some lines are weak→ need to use big telescope

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## Practical methods and examples


#### 30 Dor (HST ACS)



- Observe with spectrograph
- Measure line strengths ("fluxes") of forbidden lines and representative H recomb. line (e.g. H $\alpha)$
- Use ratios to calculate temp and density
- Calculate abundances relative to H







 $n_{\rm e}$ 

arcsec

- End of stellar existence for low- to medium- mass stars
- Important source of heavy metals which enrich ISM
- Often only source of information about abundances in external systems
- Therefore, useful to know what's in them

#### So how do we find out what's in them?

- Spectroscopic analysis
- Historically relied on brightest emission lines collisionally excited lines
- These are easily observable, but very very sensitive to physical conditions
- Small error in assumed temperature = big error in derived abundance



## **Optical Recombination Lines**

- ORLs of hydrogen and helium are strong, but those of heavy elements are much weaker
- Much less sensitive to physical conditions, so should be more reliable to derive abundances from
- Should be...

## ...but is it?

- Unfortunate problem in PNe they always give higher abundances relative to hydrogen than CELs
- Factors range from 3 30 and beyond

#### Possible solutions?

- Possibility 1: Temperature fluctuations in homogeneous gas
  - Hotter regions preferentially emit CELs, cooler regions preferentially emit ORLs. Assuming uniform temperature results in underestimated CEL abundances.

Easily ruled out with high-res spectroscopy

- Possibility 2: density fluctuations in chemically homogeneous gas
  - Very dense regions (n<sub>e</sub>~10<sup>6</sup> cm<sup>-3</sup>) would quench emission of [O III] 4959,5007 lines
  - [O III]4363 has higher critical density and is not quenched. Therefore [O III] ratio underestimated, T<sub>e</sub> overestimated, CEL abundances underestimated.

Measuring density of hydrogen gas (through method we haven't seen) rules this out

- Possibility 3: Chemical inhomogeneities
  - Strongly metal-rich zone near central stars could give rise to observed spectra – ruled out by spatially-resolved analysis
  - 2. Small metal-rich knots (too small to be directly resolved), cooled strongly by high metal content



### PN: Abell 30



## Knots: Abell 30

- Results of empirical analysis: T([O III]) ~16000 K
  T(He I) ~5000–10000 K
  T(O II) <2000 K</li>
- Abundances: ORL abundances exceed CEL abundances by factors of 300-700!!





HST imaging

## Conclusions

- Cold H-deficient knots are confirmed as existing in planetary nebulae
- Simple temperature fluctuations are ruled out from spatially resolved spectroscopy
- Evidence strongly suggests existence of cold metalrich knots in most planetary nebulae.



# **Open questions**

- No direct observations in vast majority of cases. Where did the knots come from? Why are they there?
- What is their distribution?
- What is overall chemical content of PNe?