

# Extracting the 21-cm signal\* from the cosmic dawn

\*(especially the sky-averaged 21-cm signal)

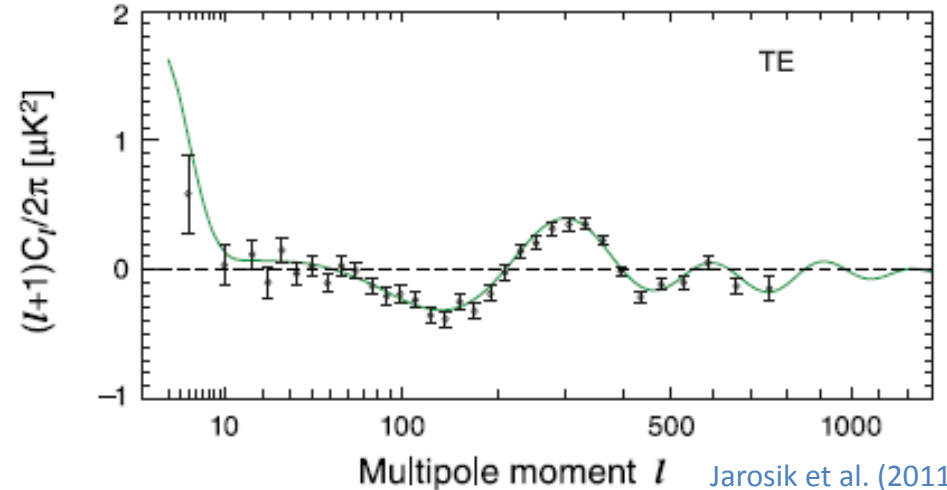
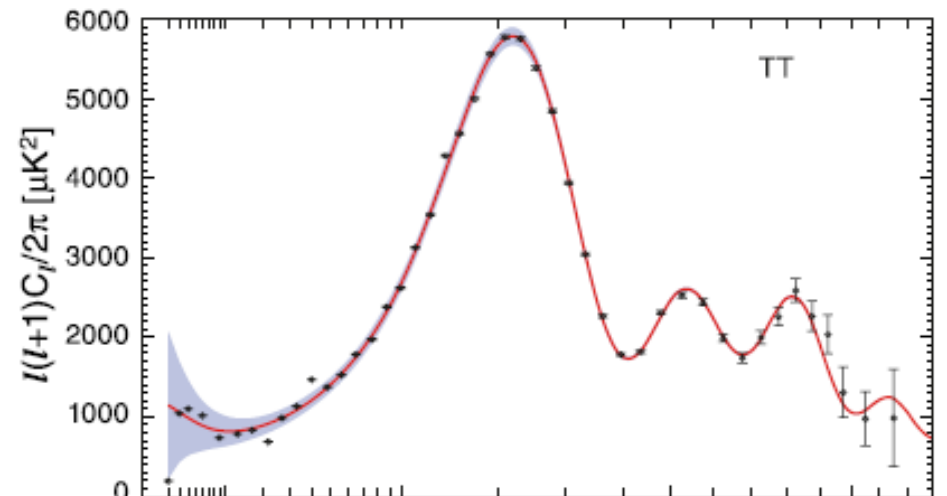
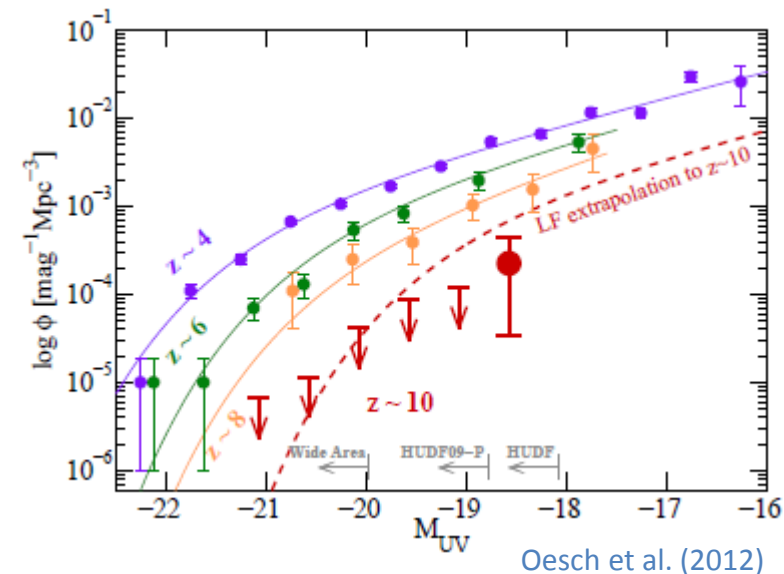
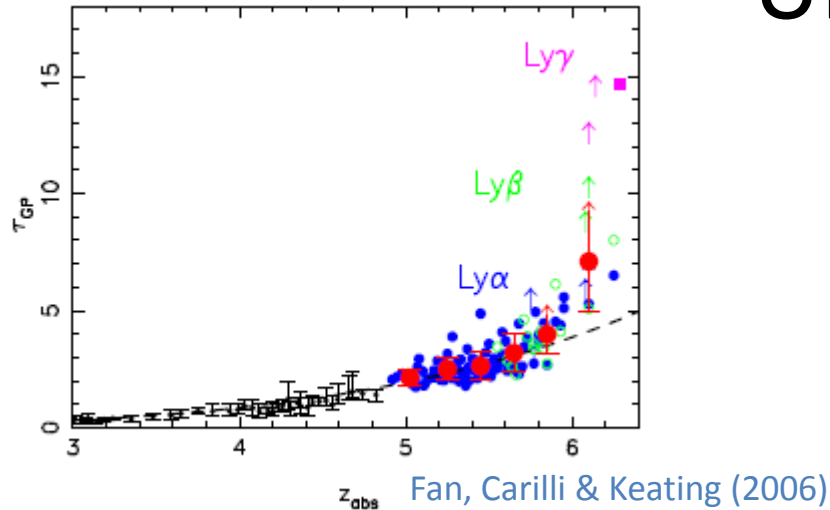
Geraint Harker

# Stages in the evolution of the Universe, from redshift 1100 to 6



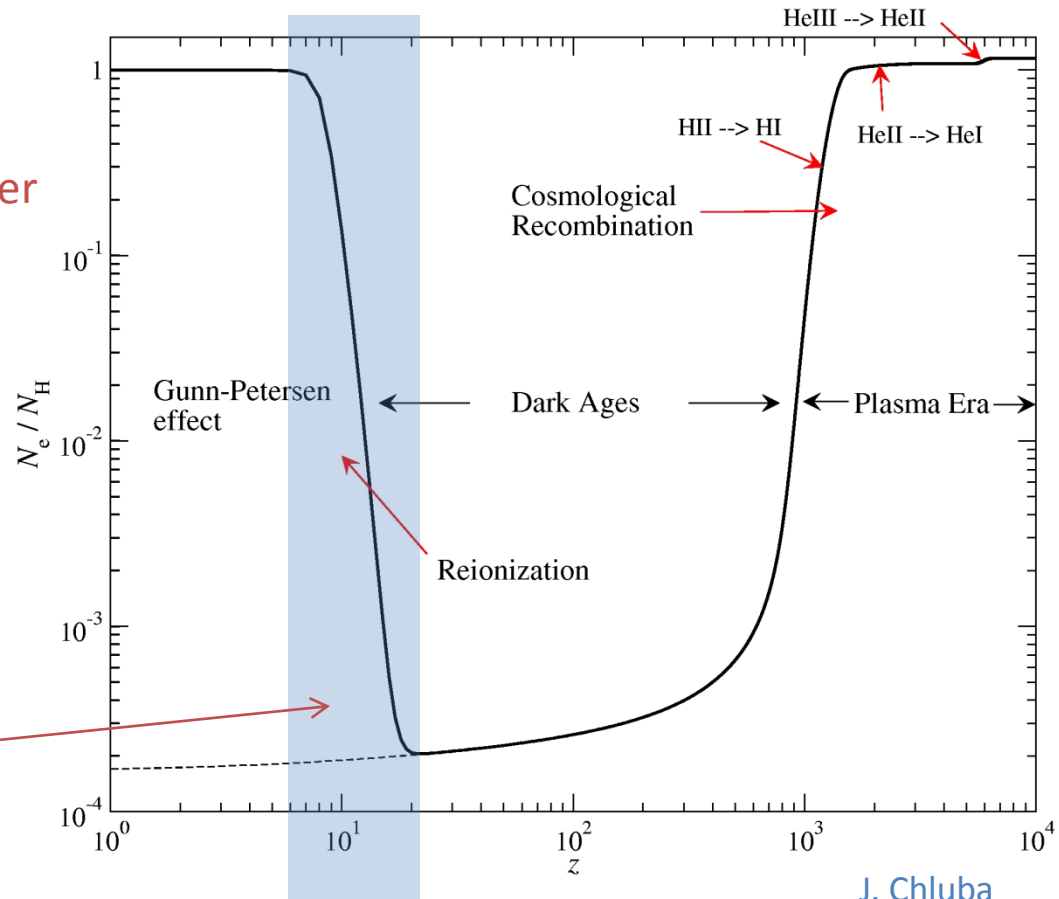
DARE proposal; graphics adapted from A. Loeb, 2006, *Scientific American*, 295, 46

# Observational constraints on the $z > 6$ Universe



# Ionization fraction

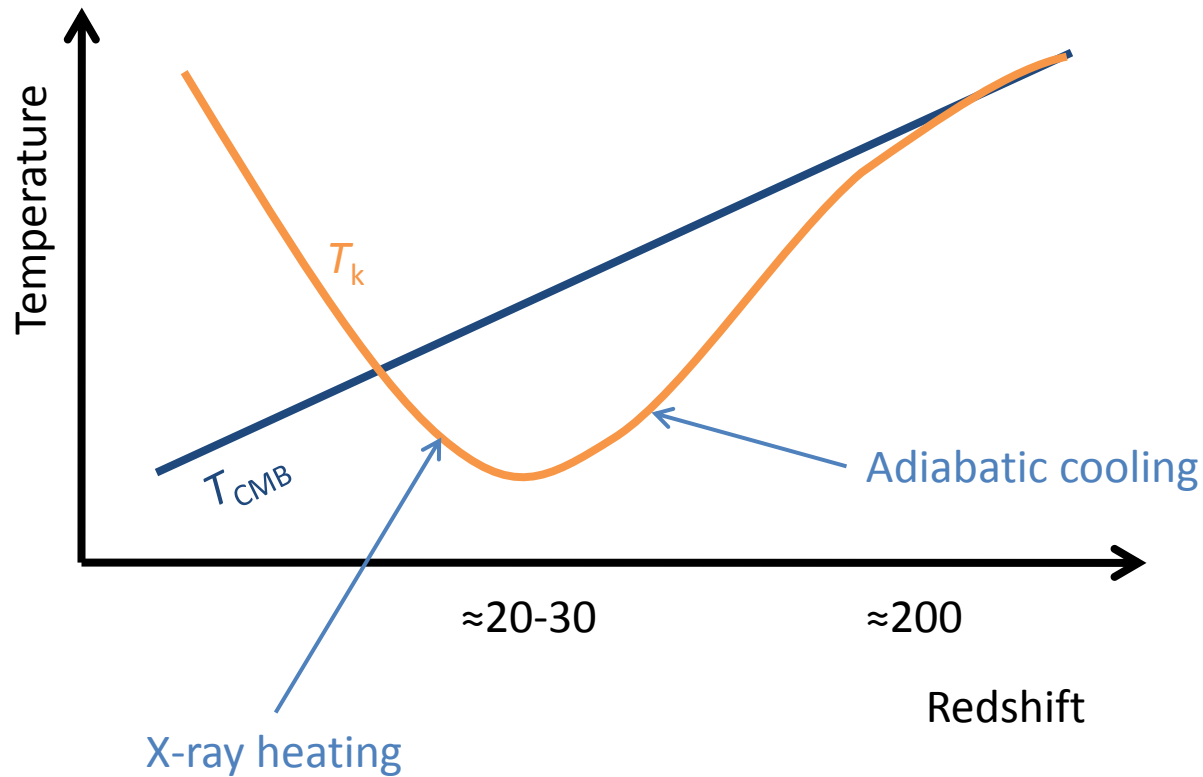
Number density of free electrons divided by number density of hydrogen atoms



Epoch of reionization (EoR)

J. Chluba

# Temperature evolution



# The hydrogen 21-cm line

$$\frac{n_1}{n_0} = \frac{g_1}{g_0} e^{-T_*/T_{\text{spin}}} \quad (T_* = 0.068\text{K})$$

'Colour temperature', shown to be equal to  $T_k$  for the situation studied here

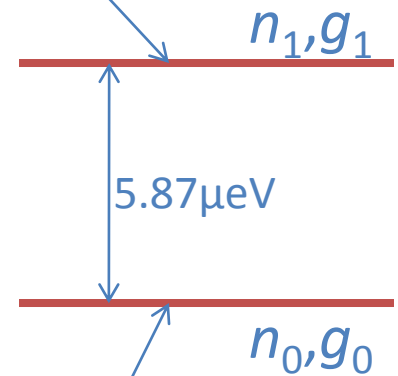
$$T_{\text{spin}} = \frac{T_{\text{CMB}} + y_\alpha T_\alpha + y_c T_k}{1 + y_\alpha + y_c}$$

Field, 1958, Proc. IRE

$$\delta T_b = \frac{T_{\text{spin}} - T_{\text{CMB}}}{1 + z} \underbrace{(1 - e^{-\tau_{\nu 0}})}_{\approx \tau_{\nu 0} \propto 1/T_{\text{spin}}}$$

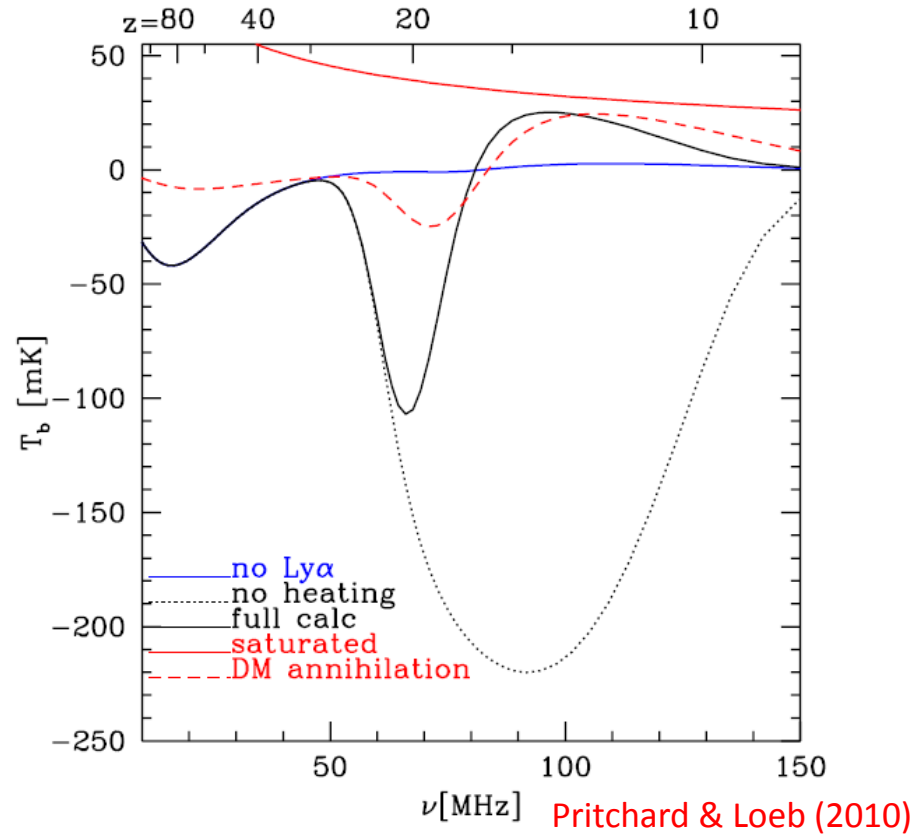
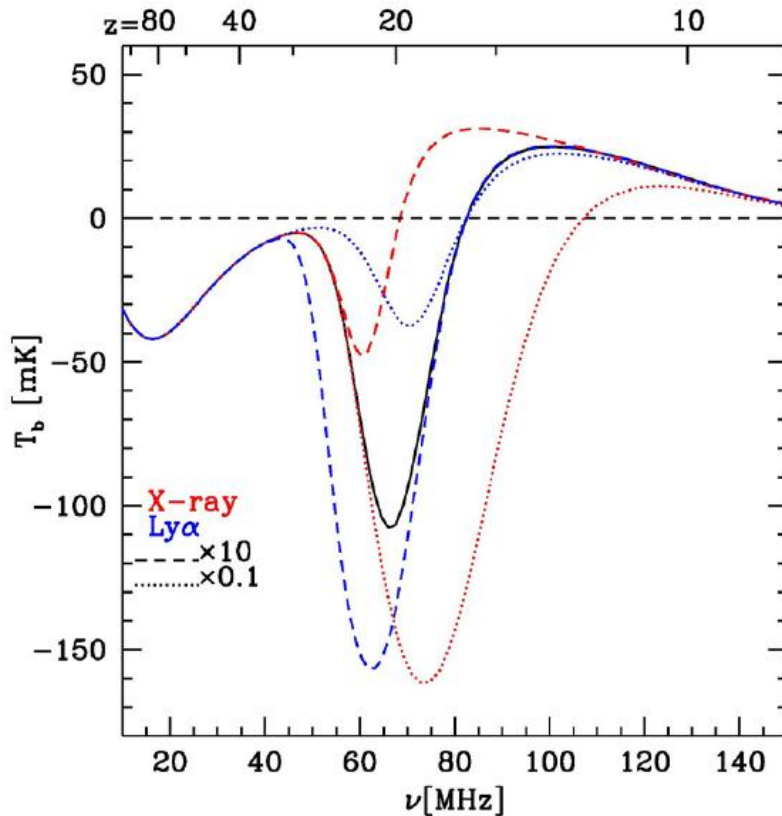
Hydrogen 1s ground state

Parallel spins



Antiparallel spins

# The global 21-cm signal

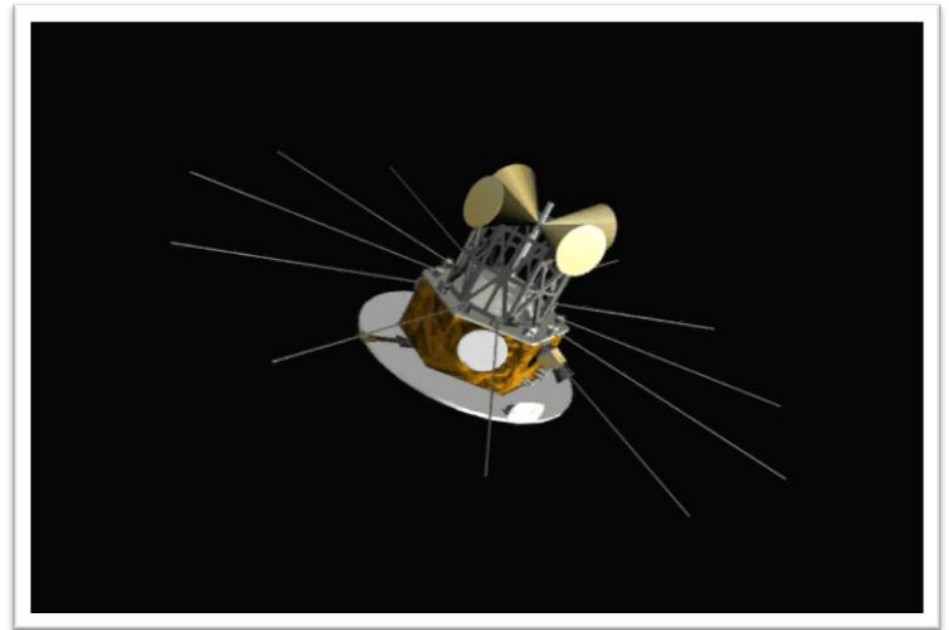


Pritchard & Loeb (2010)



# Global 21-cm experiments

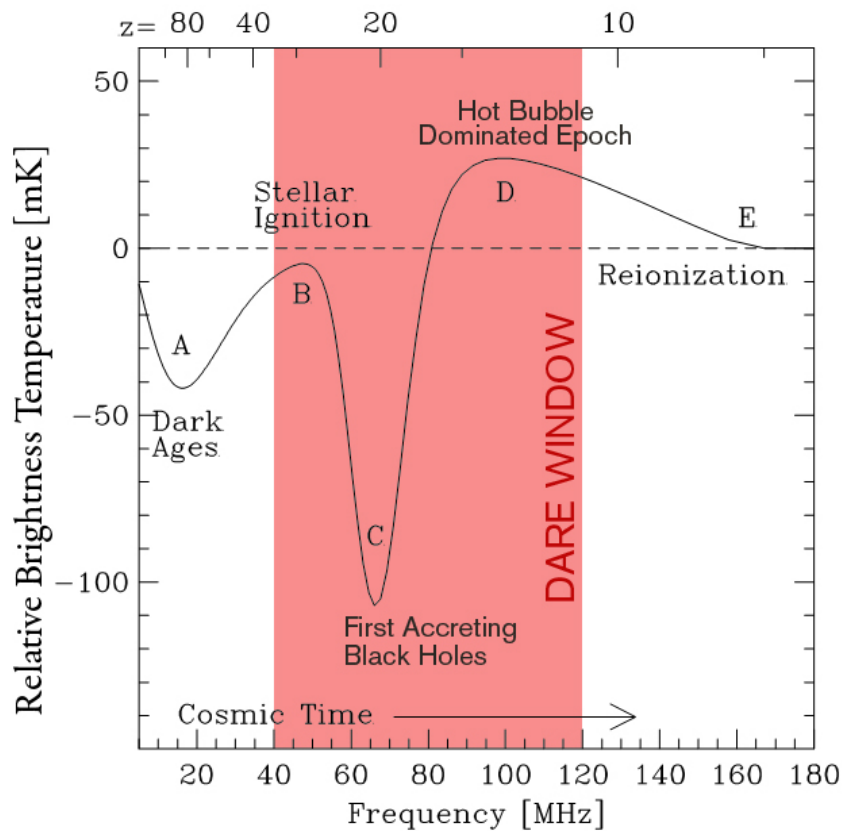
- DARE
- EDGES
- CoRE/CoRE2
- BIGHORNS
- LEDA (LWA)



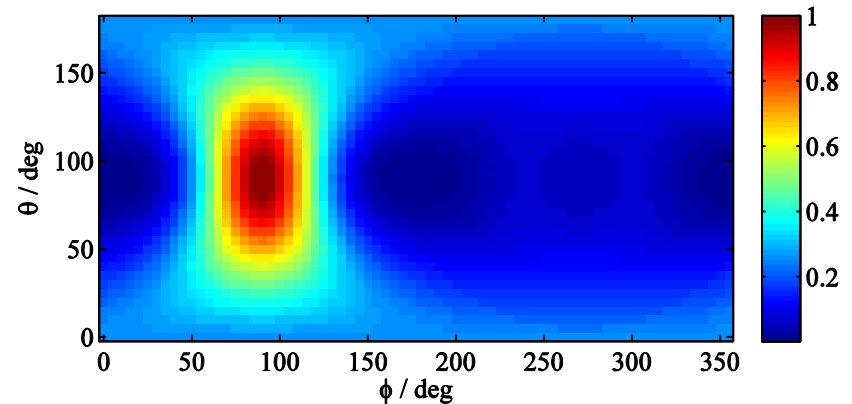
- Operates over the lunar farside
- Escapes RFI
- Whole sky available; beam covers  $\approx 1/8$  of the sky
- No ionospheric distortion or contribution to the spectrum.



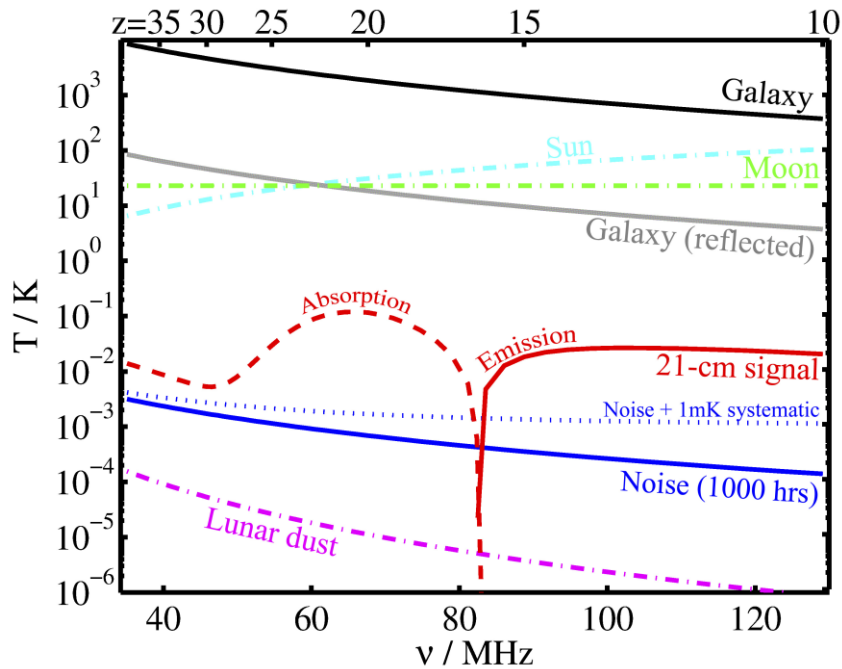
# Basic parameters of the DARE experiment



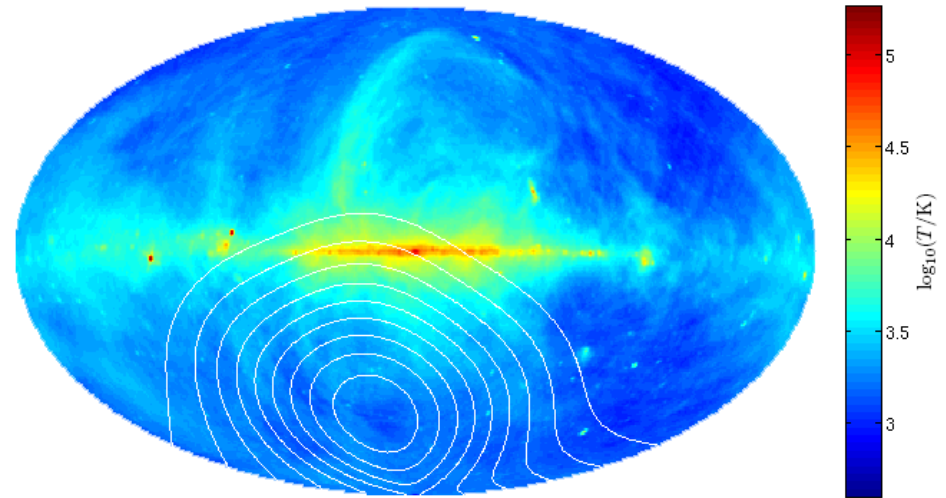
DARE antenna power pattern at 75 MHz



# Foregrounds



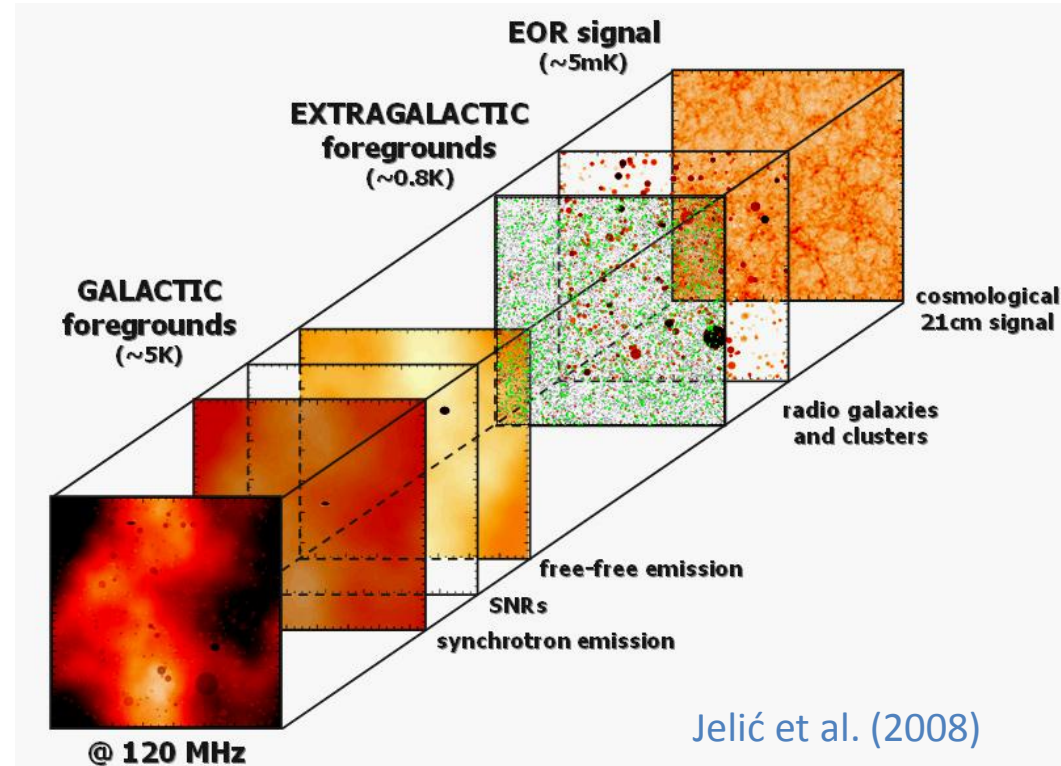
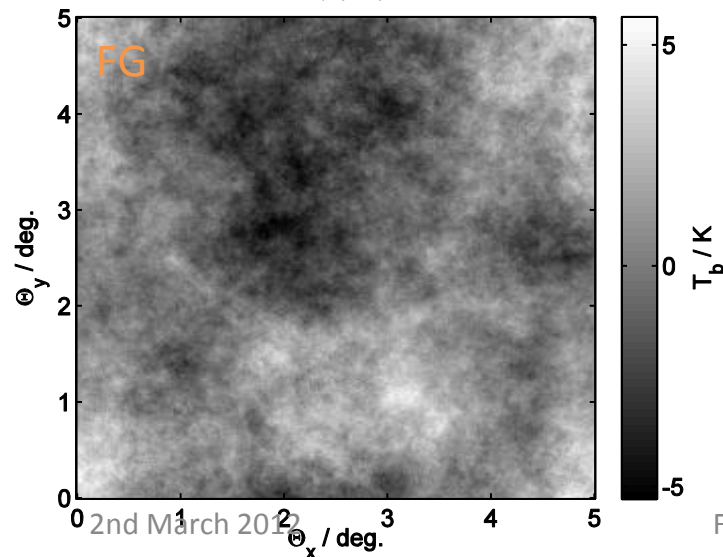
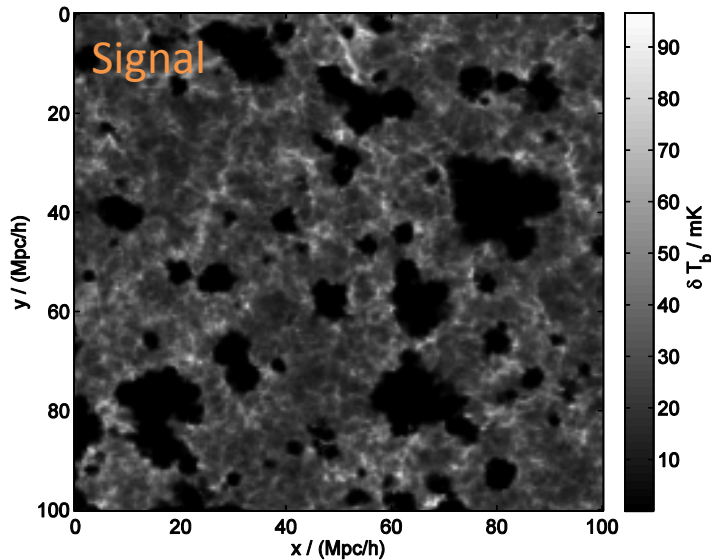
Spectrally smooth...



de Oliveira-Costa et al. (2008)

... but spatially variable

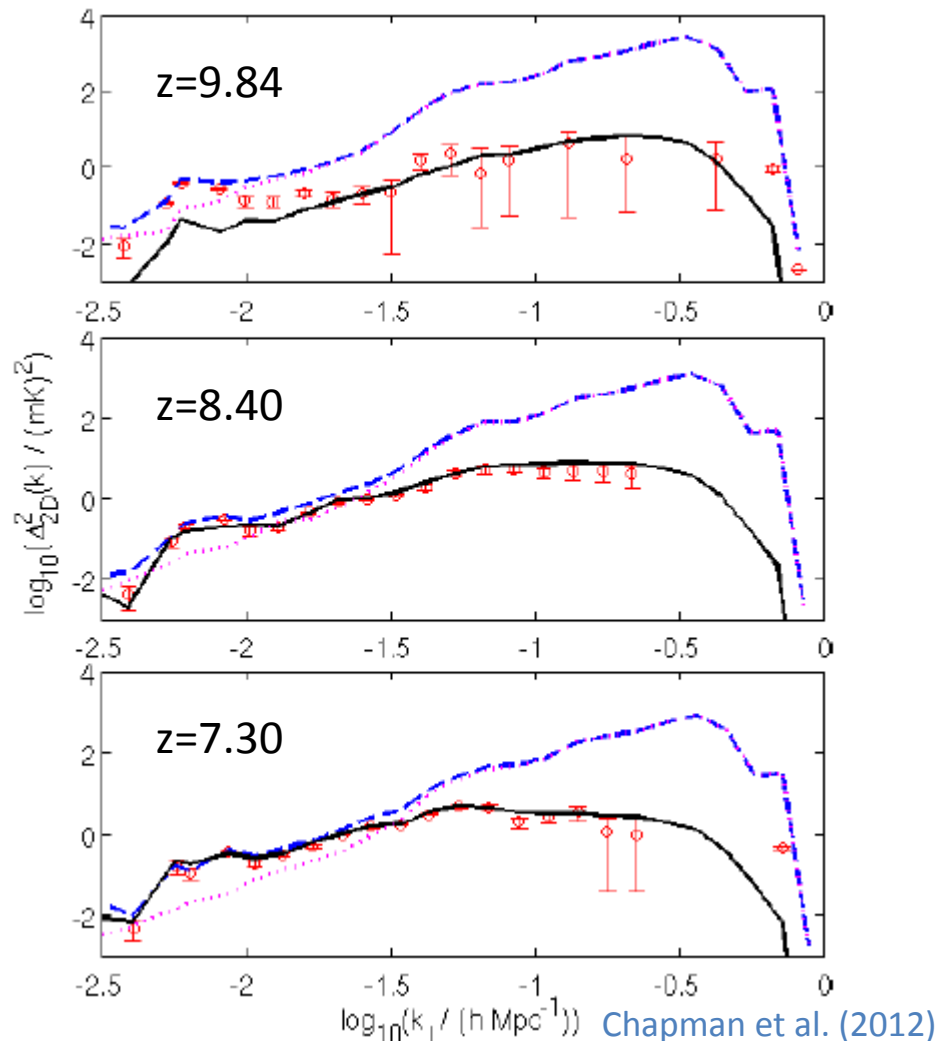
# Similar-looking problem for experiments going after spatial information



Jelić et al. (2008)

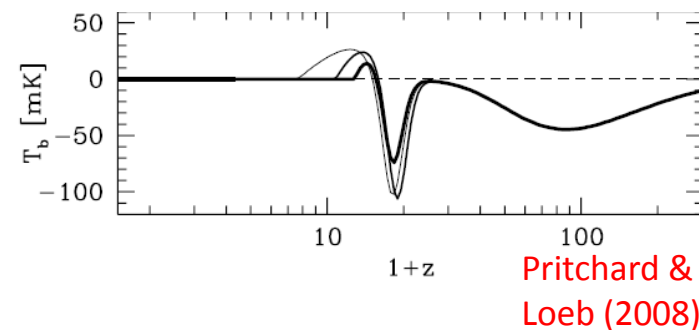
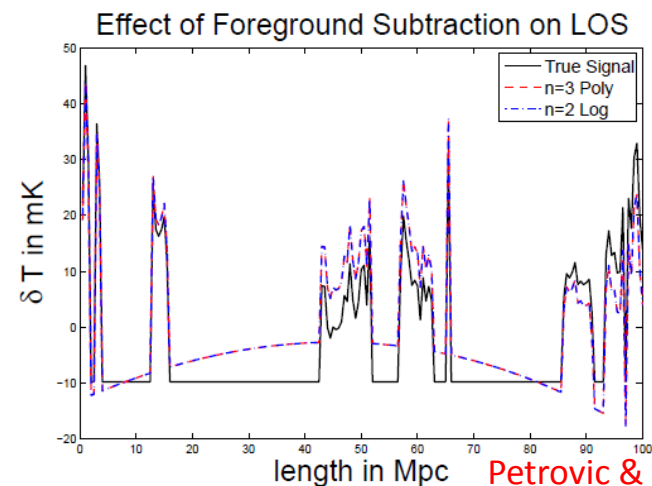
# Independent component analysis

- Sophisticated foreground removal tools have been developed for the CMB and for interferometric 21-cm experiments.
- ICA uses the statistical independence of the foregrounds and signal to separate the two with minimal assumptions.
- Uses the fact that we have many different samples of the signal and foreground from the many pixels of the interferometric map.
- Other approaches:
  - Wp smoothing: uses smoothness of foregrounds directly
  - Even fitting simple parametric models seems to work OK.



# Interferometric and sky-averaged 21-cm foregrounds: similarities

- Foregrounds dominate over signal by orders of magnitude, wherever you look in the sky.
- Use the different spectral structure of the foregrounds and 21-cm signal to distinguish them: there are good reasons to think that many of the foregrounds are spectrally smooth.
- The spatial correlation of signal and foregrounds are also different, though this is less often exploited.
- The foregrounds and the instrument are coupled together strongly: can't remove the foregrounds without understanding both (c.f. simultaneous fitting of signal and instrument in e.g. FIRAS analysis).



# Interferometric and sky-averaged 21-cm foregrounds: differences

- Averaged over a big enough area of sky, the global signal is the same wherever you look whereas the foregrounds vary (could help with subtraction, as suggested by Shaver et al. 1999).
- Point sources are dealt with very differently:
  - Carefully subtracted for interferometer experiments, but see e.g. Datta, Bowman & Carilli (2010) or Morales et al. (2012).
  - Averaged over and treated as a diffuse foreground for global signal experiments.
- The signal is a lot smoother in the sky-averaged case, and the foregrounds are effectively much larger (especially for ‘cosmic dawn’ / ‘dark ages’ work), so stronger assumptions need to be made about the foregrounds (and calibration of the frequency response becomes even more crucial: we want a stable environment!).
- Easier to beat down the noise below the level of the signal for the global signal.
- Nothing to cross-correlate the global signal with?
- RFI for the global signal could be even more awkward: can’t be localised, may require a more complicated receiver design, etc., though a single antenna experiment could in principle be much simpler and cheaper.



# Recovering the shape of the global 21-cm signal from simulated DARE data

Developed parametrized models of the signal and foregrounds in eight directions:

- Galaxy and diffuse extragalactic sources
- Sun
- Moon (emission and reflections)
- Instrument



- Simulate data
- Fit the parameters and derive errors with a Markov Chain Monte Carlo code

# Bayesian inference

Parameters

Data

Posterior

Hypothesis or model

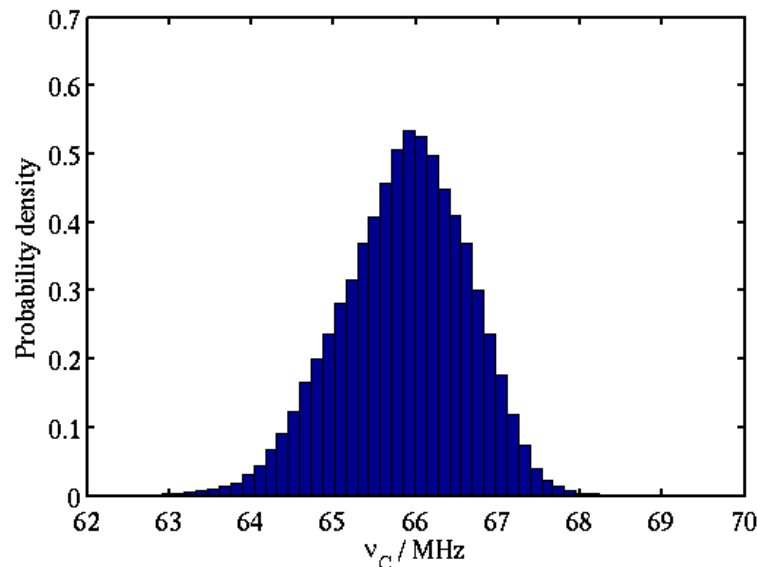
Likelihood,  $\mathcal{L}(\Theta)$

Prior,  $\pi(\Theta)$

Evidence,  $\mathcal{Z}$

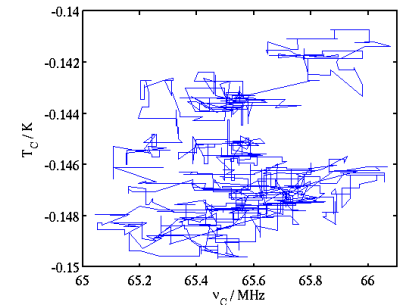
$$\Pr(\Theta | \mathbf{D}, H) = \frac{\Pr(\mathbf{D} | \Theta, H) \Pr(\Theta | H)}{\Pr(\mathbf{D} | H)}$$
$$\mathcal{Z} = \int \pi(\Theta) \mathcal{L}(\Theta) d^N \Theta$$
The diagram illustrates the Bayesian inference equation. On the left, the posterior probability  $\Pr(\Theta | \mathbf{D}, H)$  is enclosed in an orange box. An arrow labeled 'Parameters' points to  $\Theta$ , and another arrow labeled 'Data' points to  $\mathbf{D}$ . Below the box is the label 'Posterior'. An arrow labeled 'Hypothesis or model' points to  $H$ . The equation is followed by an equals sign and a fraction. The numerator consists of two terms:  $\Pr(\mathbf{D} | \Theta, H)$  enclosed in a purple box and labeled 'Likelihood,  $\mathcal{L}(\Theta)$ ', and  $\Pr(\Theta | H)$  enclosed in a green box and labeled 'Prior,  $\pi(\Theta)$ '. The denominator is  $\Pr(\mathbf{D} | H)$  enclosed in a blue box and labeled 'Evidence,  $\mathcal{Z}$ '. Below the entire equation, the integral formula for the evidence is given:  $\mathcal{Z} = \int \pi(\Theta) \mathcal{L}(\Theta) d^N \Theta$ .

# The Markov Chain Monte Carlo technique



A Markov Chain Monte Carlo simulation allows us to draw unbiased, random samples from the posterior probability distribution of the parameters we're trying to find.

The path taken by part of the Markov Chain through a two-dimensional slice of parameter space. The parameter space has 73 dimensions in our model.

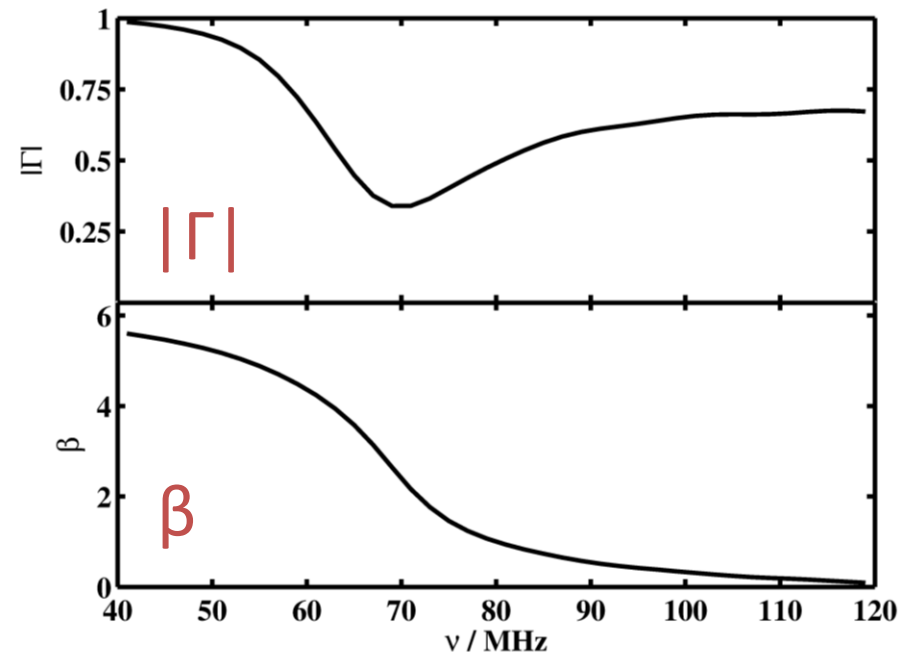
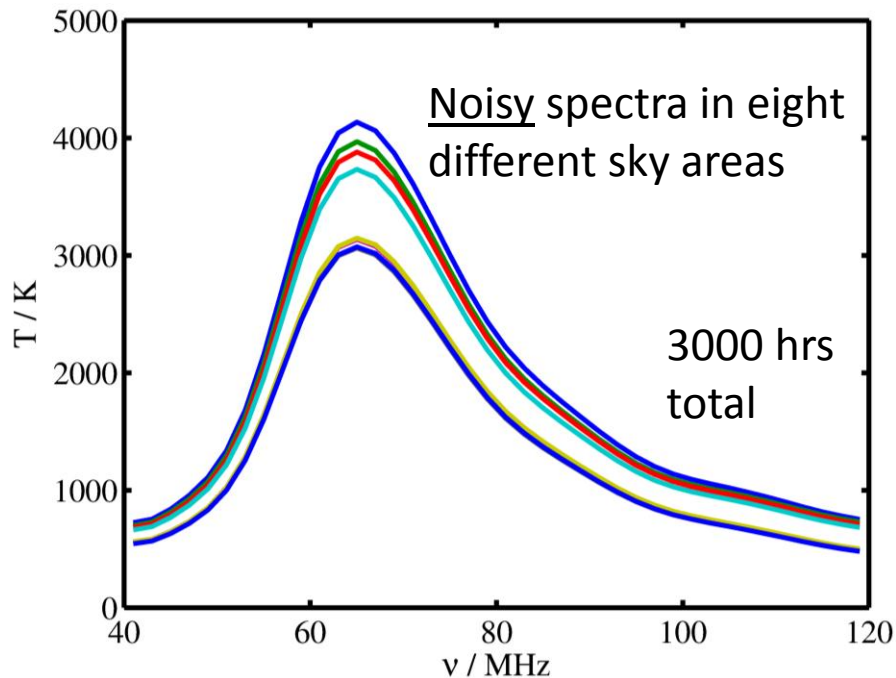


Parameter group	No. of parameters
21-cm signal	$3 \times 2 = 6$
Diffuse foregrounds	$4 \times 8 = 32$
Sun	$8 + 3 = 11$
Moon	2
Instrument	22
<b>Total</b>	<b>73</b>

# Instrument frequency response and simulated spectra

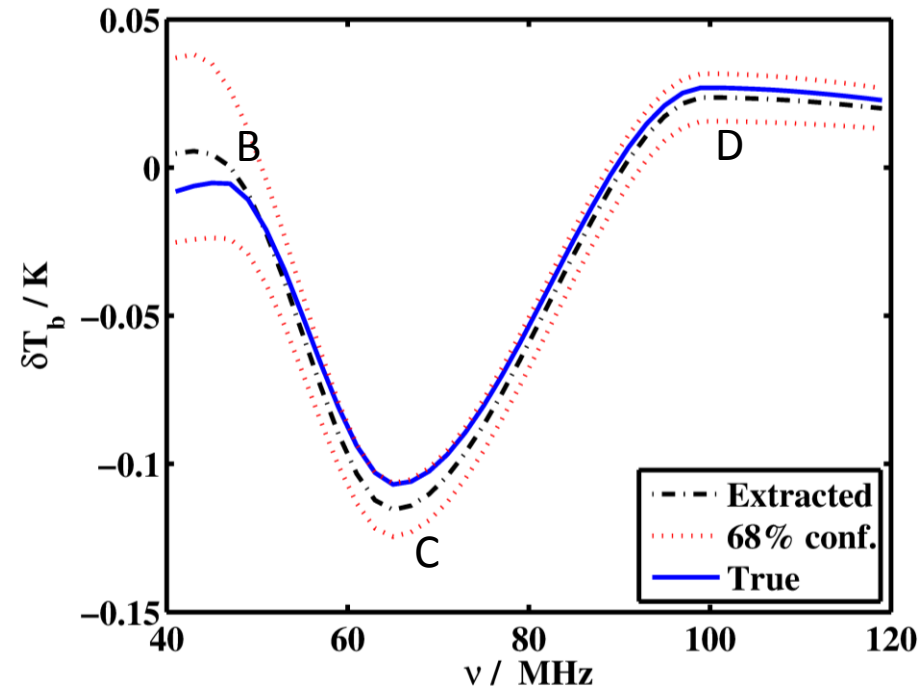
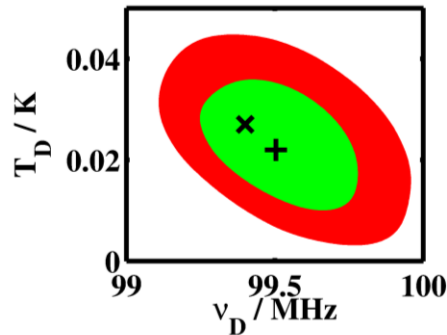
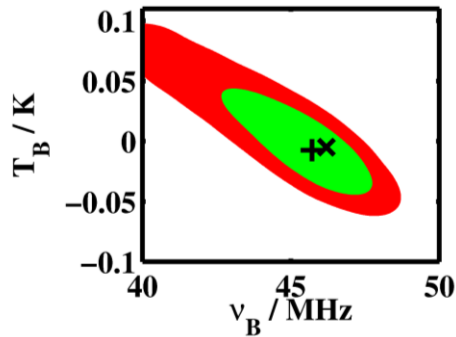
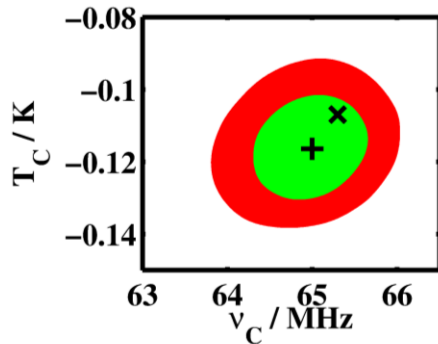
$$T_{\text{ant}}(\nu) = T_a + |\Gamma(\nu)|^2 T_b + 2T_c |\Gamma(\nu)| \cos[\beta(\nu) + \phi_c] + T_{\text{sky}}(\nu) [1 - |\Gamma(\nu)|^2]$$

Take  $T_c = \epsilon T_b = \epsilon T_a$

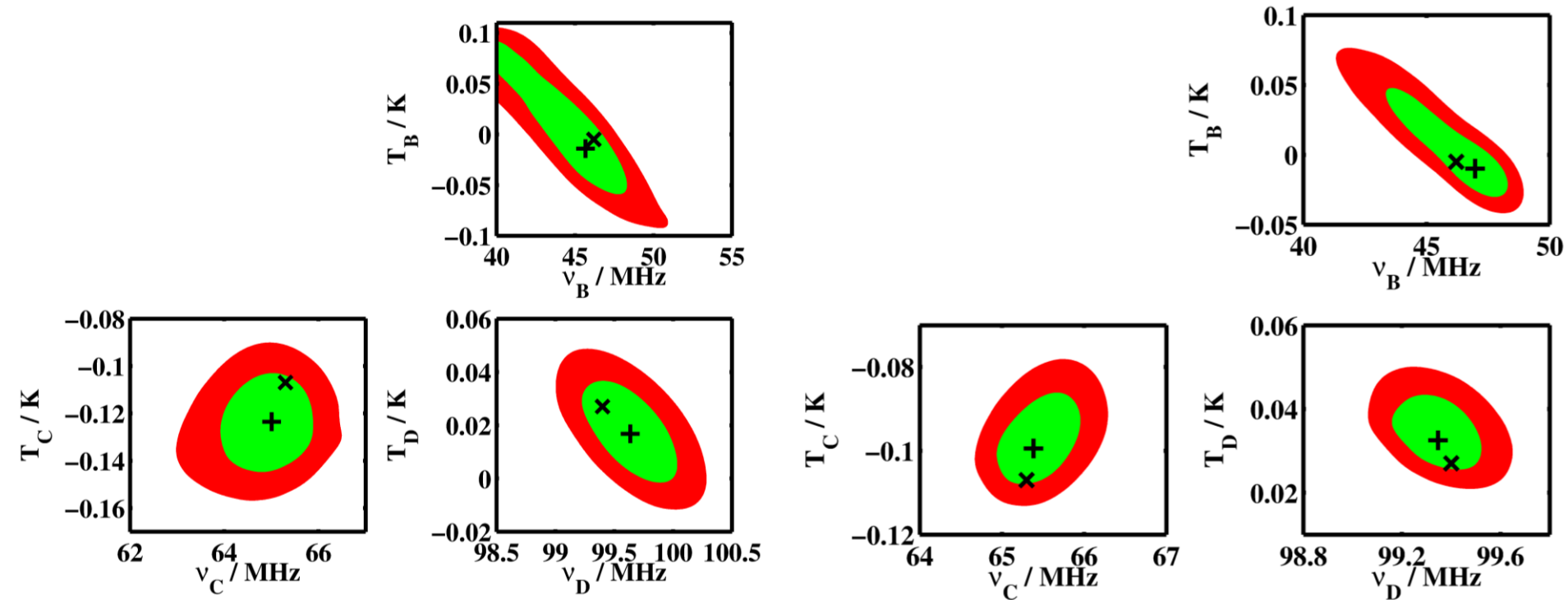


# MCMC results: positions of turning points and shape of signal (3000 hrs)

x Input  
+ Recovered  
• 68% conf.  
• 95% conf.



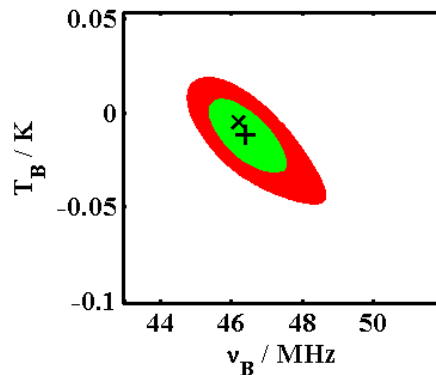
# MCMC results: 1000 and 10000 hours



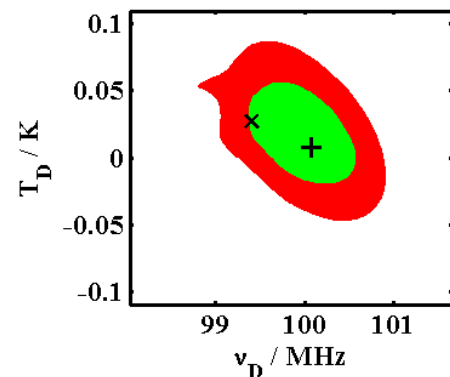
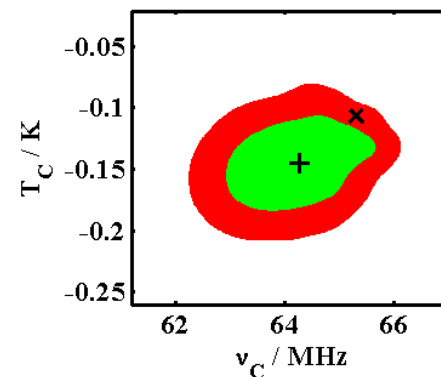
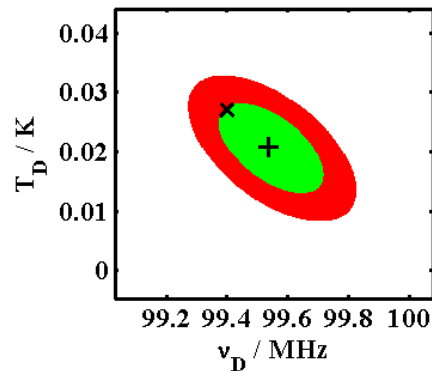
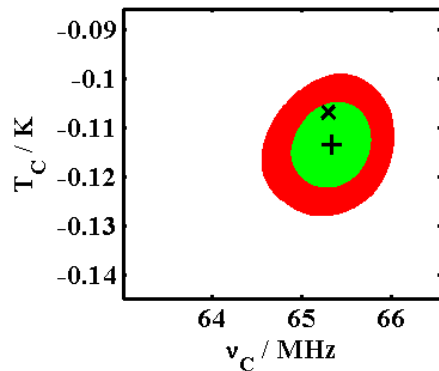
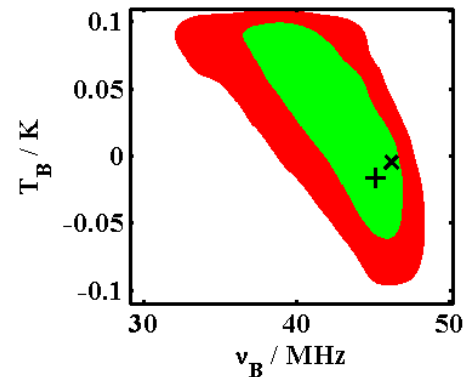


# Varying the assumptions

Redesigned  
antenna (wider  
cone opening  
angle)

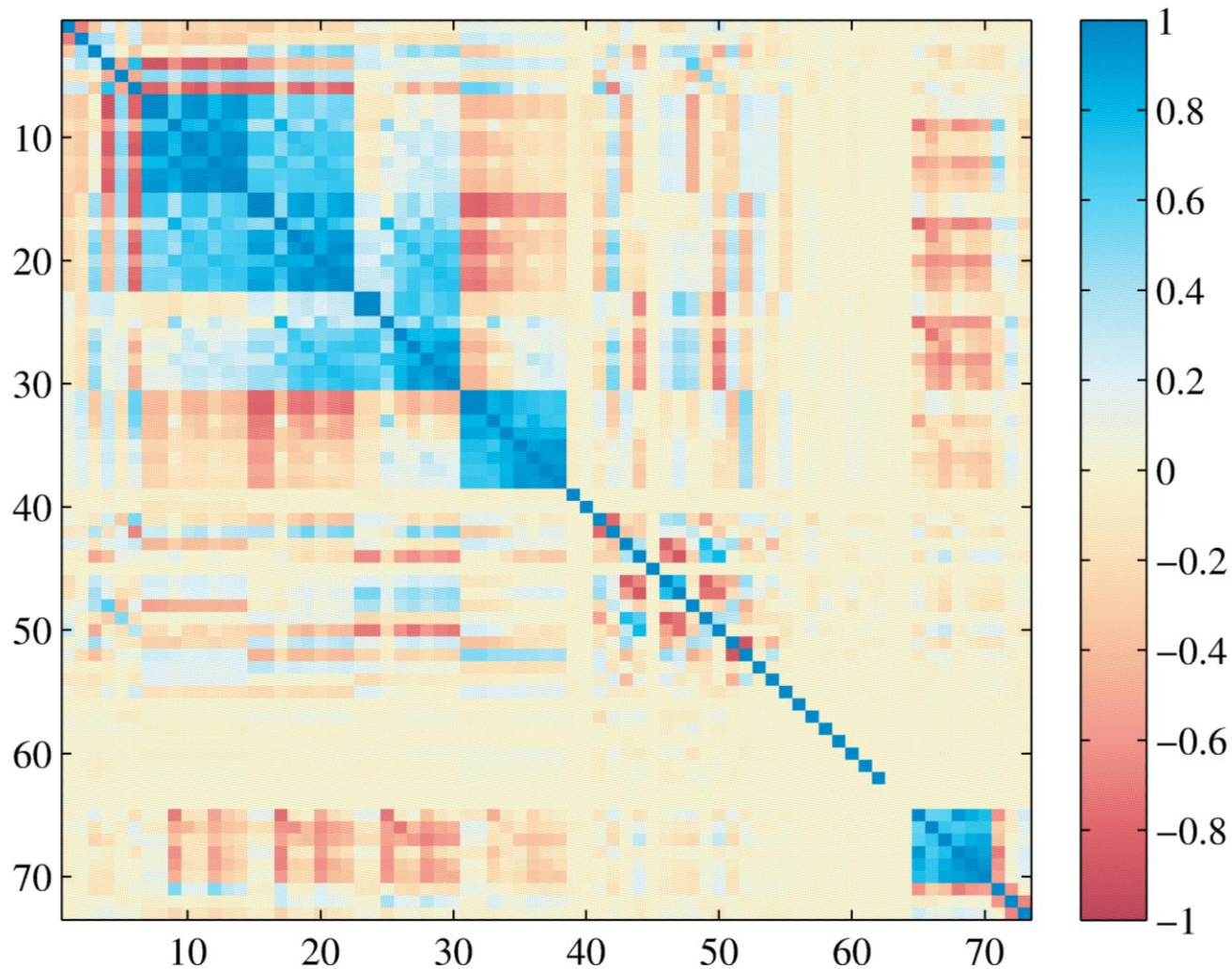


Looser priors  
on instrument  
response



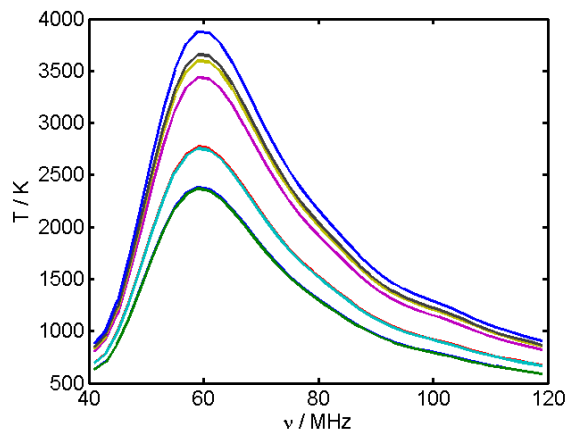
Constraints have actually been improved a little with a slightly modified antenna design!  
But accurate modelling of the instrument remains the crucial ingredient.

# Correlation matrix

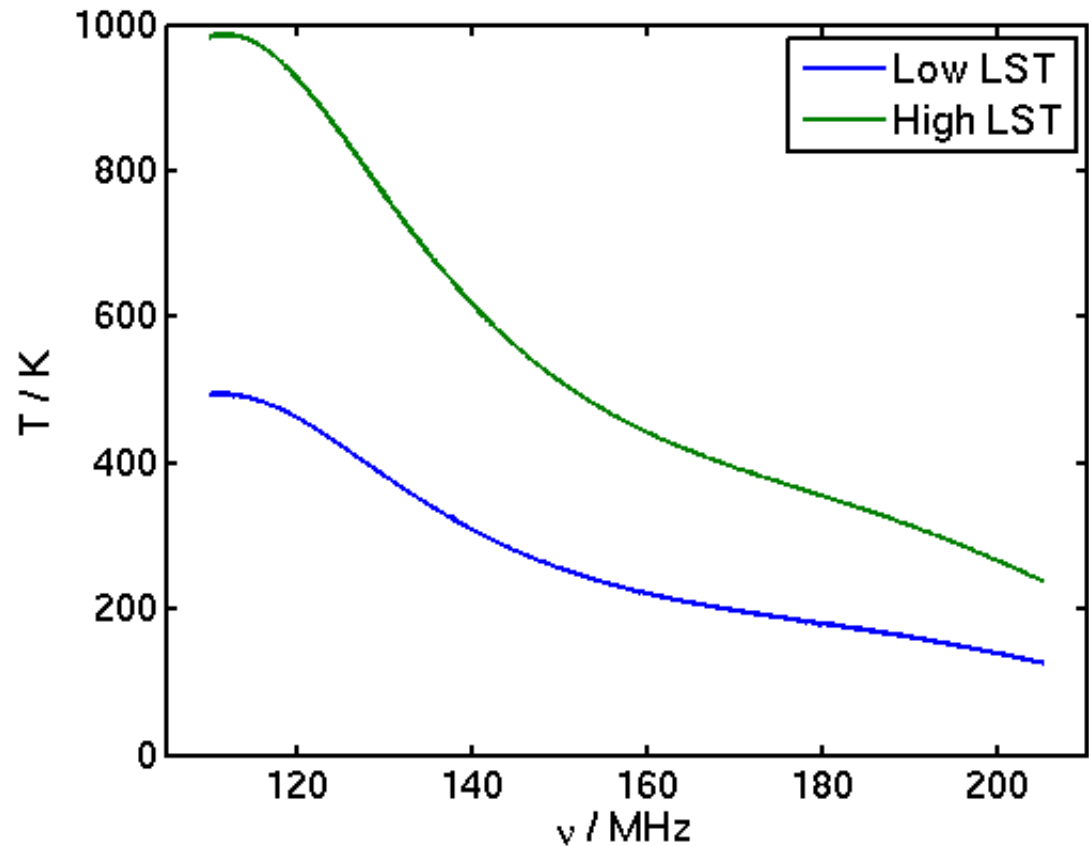


# Applying MCMC to EDGES data

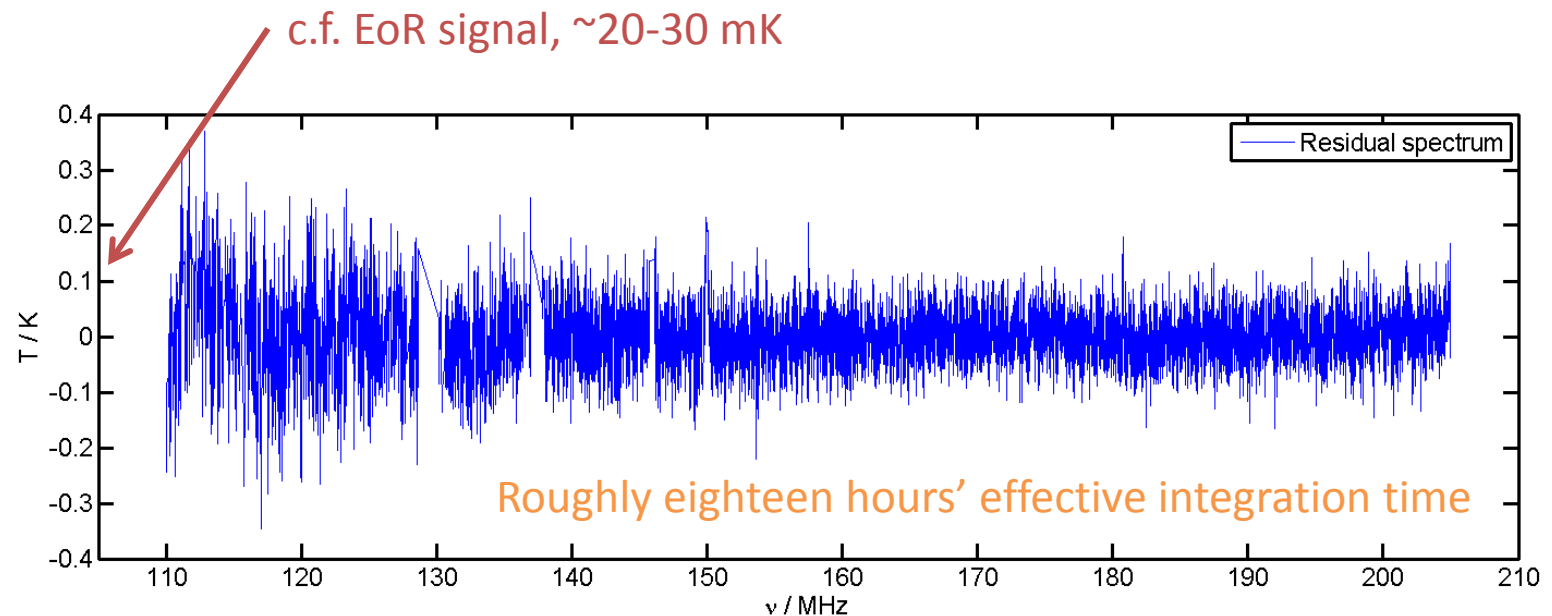
How does this analysis pipeline work with real data? Can we improve constraints on the epoch of reionization data?



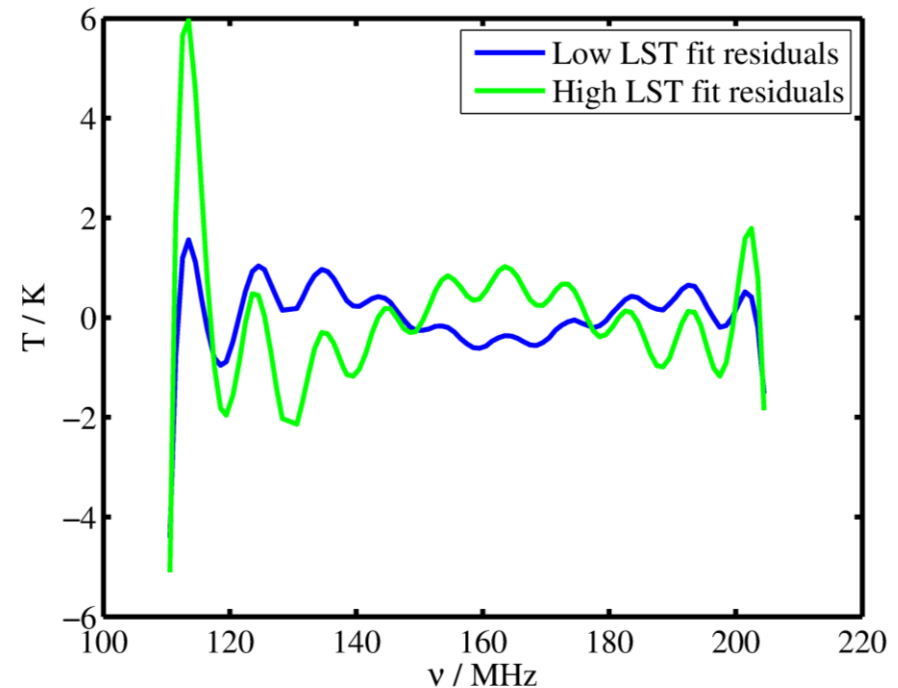
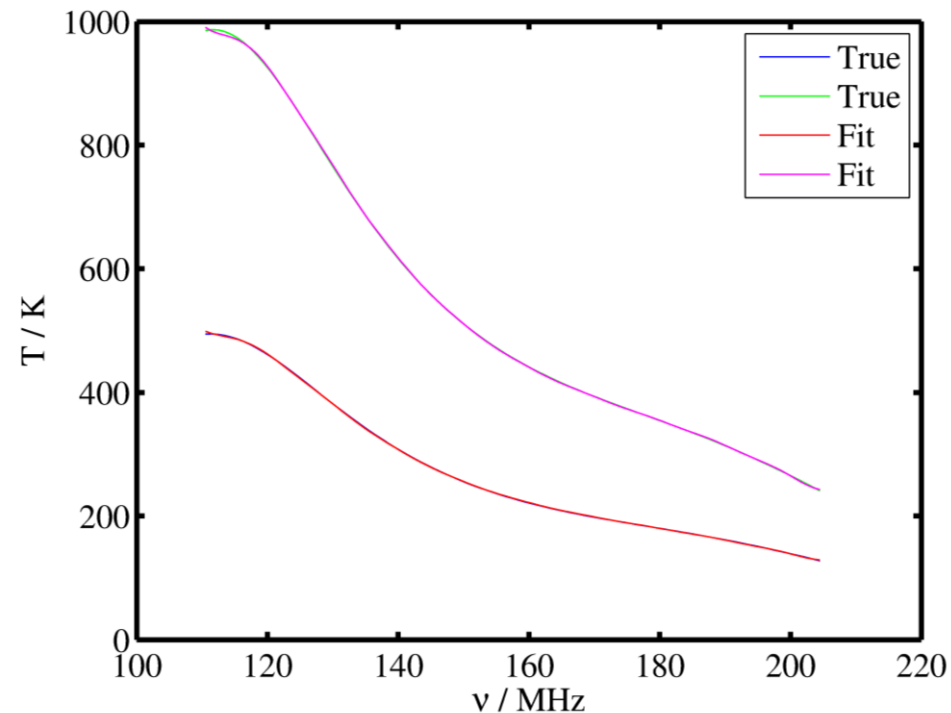
c.f. simulated  
DARE data



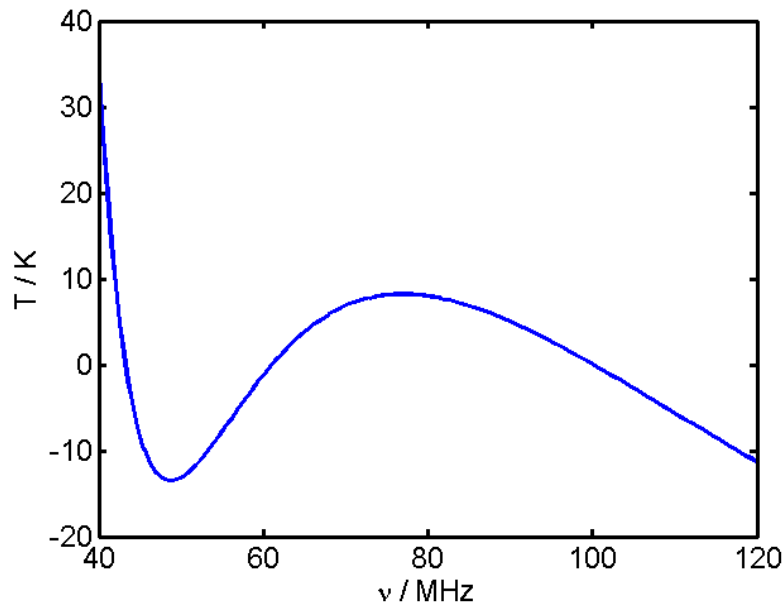
# Applying MCMC to EDGES data



# Applying MCMC to EDGES data



# The ionosphere in global signal experiments



Fitting residuals for a power-law galaxy spectrum and an unmodelled ionospheric contribution

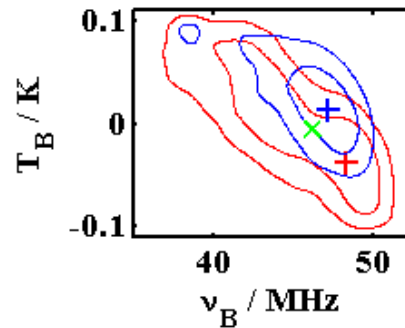
- The ionosphere absorbs low-frequency radio, and contains a population of  $\sim 1000$  K electrons which also radiate at low frequencies.
- Importance of these effects goes like  $(\text{frequency})^{-2}$
- Absorbs  $\sim 1\%$  of radiation at these frequencies.
- Simple error analysis (A.E.E. Rogers, EDGES memo 79) finds that reionization parameter errors are roughly doubled.
- Adds another 2 parameters per sky region to our model.

$$T_{\text{with}} = T_{\text{without}} - T_{\text{without}}(1 - L) \left( \frac{\nu}{\nu_0} \right)^{-2} + T_{\text{electrons}}(1 - L) \left( \frac{\nu}{\nu_0} \right)^{-2}$$

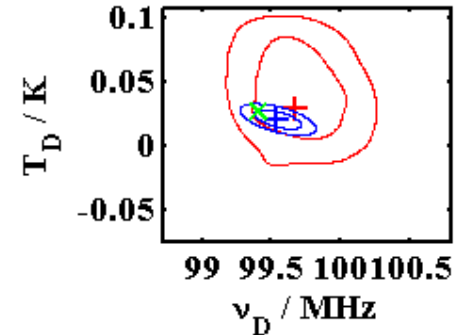
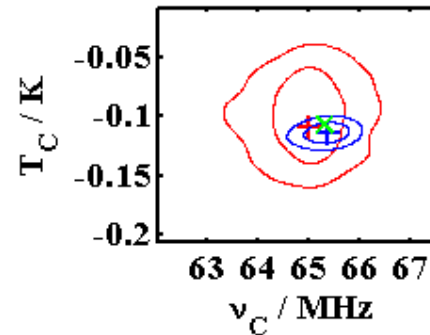
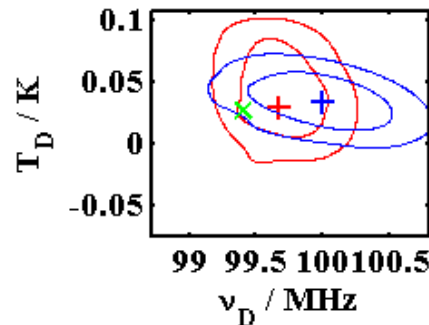
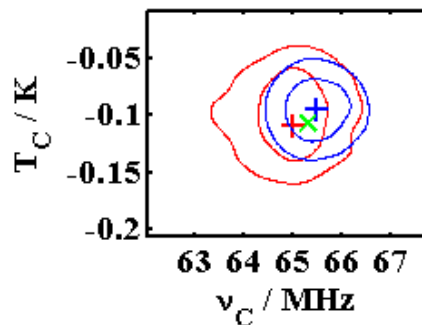
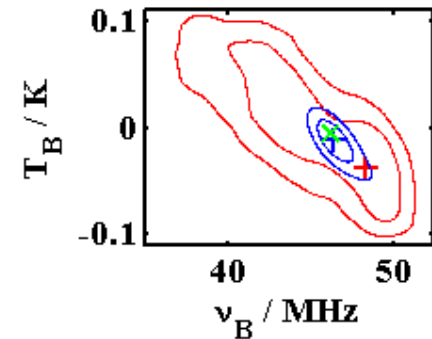


# (Fundamental?) limits to a ground-based experiment

Loose priors on space-based instrument



Tight priors on space-based instrument

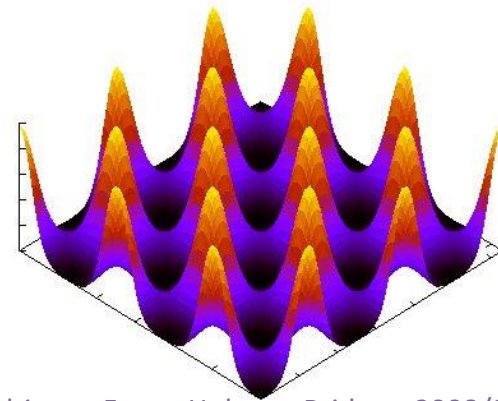


Blue = space measurement, 8 sky regions, 3000 hours

Red = ground-based measurement, 2 sky regions, 10000 hours

# Coming up: code development and the DARE prototype system

- Increase the power and flexibility of the MCMC code:
  - Incorporate existing code base developed by other groups.
  - Find a way to start from high-resolution, time-ordered satellite data rather than assuming we begin with preprocessed data.
  - Include a wider range of 21-cm models: do *model selection* rather than simply parameter estimation.



Multinest: Feroz, Hobson, Bridges, 2008/9

# Bayesian inference

Parameters

Data

$\Pr(\Theta | \mathbf{D}, H)$  =  $\frac{\Pr(\mathbf{D} | \Theta, H) \Pr(\Theta | H)}{\Pr(\mathbf{D} | H)}$

Posterior

Likelihood,  $\mathcal{L}(\Theta)$

Prior,  $\pi(\Theta)$

Evidence,  $\mathcal{Z}$

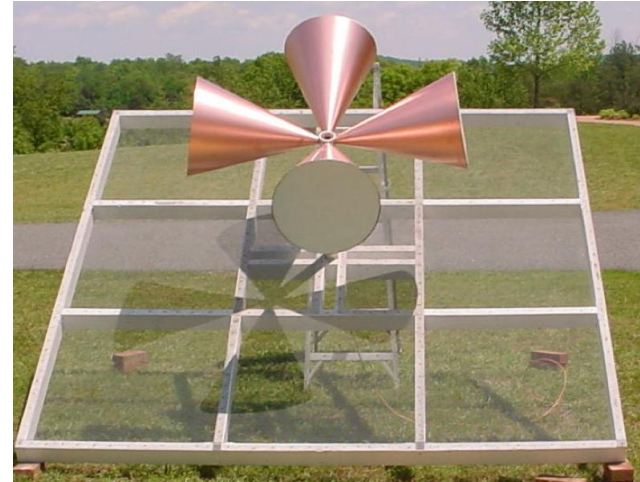
Hypothesis or model

The diagram illustrates Bayes' theorem. On the left, the posterior probability  $\Pr(\Theta | \mathbf{D}, H)$  is enclosed in an orange box. An arrow labeled 'Parameters' points to  $\Theta$ , and another arrow labeled 'Data' points to  $\mathbf{D}$ . Below the box is the label 'Posterior'. An arrow labeled 'Hypothesis or model' points to  $H$ . To the right of the equals sign, the numerator consists of two terms:  $\Pr(\mathbf{D} | \Theta, H)$  enclosed in a purple box and labeled 'Likelihood,  $\mathcal{L}(\Theta)$ ', and  $\Pr(\Theta | H)$  enclosed in a green box and labeled 'Prior,  $\pi(\Theta)$ '. The denominator is  $\Pr(\mathbf{D} | H)$  enclosed in a blue box and labeled 'Evidence,  $\mathcal{Z}$ '.

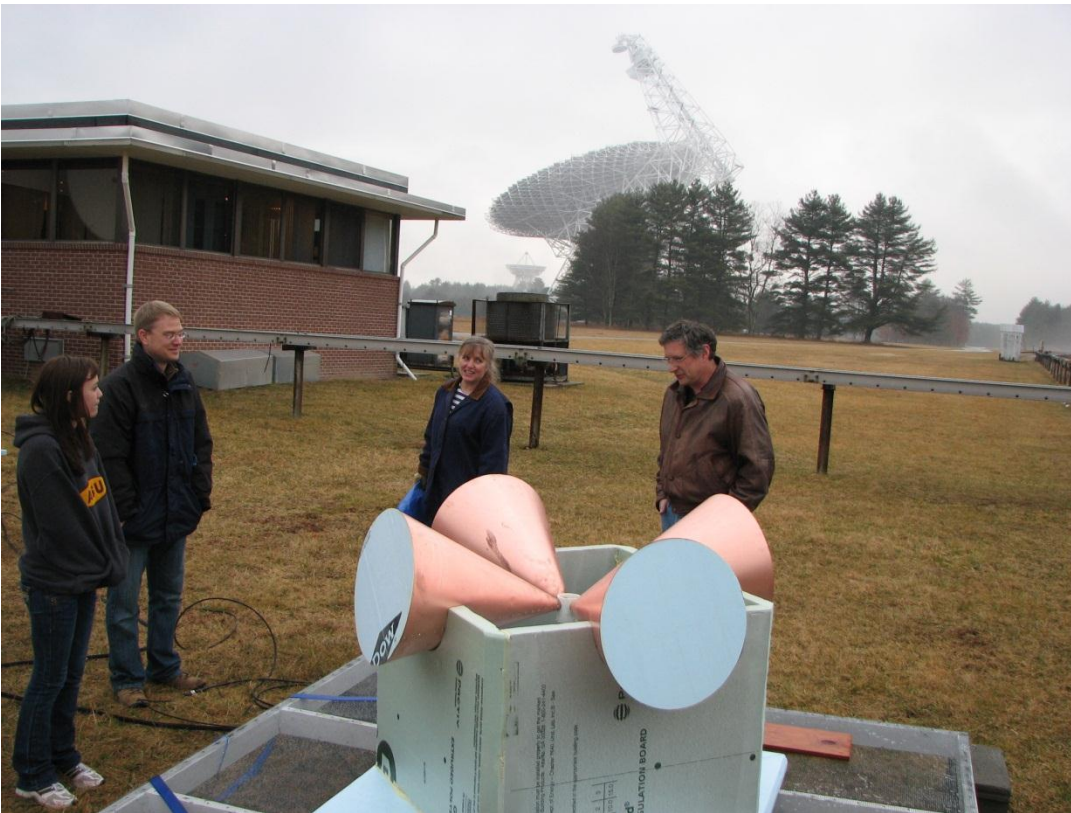
$$\mathcal{Z} = \int \pi(\Theta) \mathcal{L}(\Theta) d^N \Theta$$

# Coming up: code development and the DARE prototype system

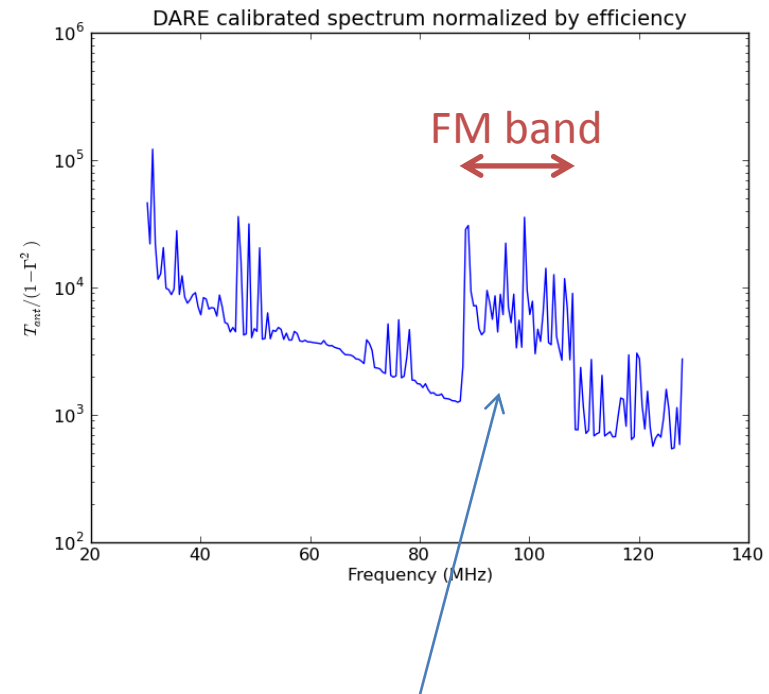
- Applying the MCMC code to data from the DARE prototype system will be a good test of the code and will also require further development:
  - Incorporating the effects of environmental changes, solar bursts etc. will require the use of the time-ordered data
  - Does the ionospheric modelling in the current version of the MCMC code hold up?
  - Are tight constraints on the 21-cm signal possible using this or EDGES? Can we prove it?



# The DARE prototype is up and running!



Some telescopes at Green Bank (pic and spectrum courtesy Abhi Datta)



Hopefully looks nicer from Western Australia (where the prototype gets shipped this month)

# Summary

- Although foreground subtraction for sky-averaged experiments shares some features with interferometric experiments, it is different enough that we need different techniques.
- The foregrounds, signal and instrumental properties probably need to be measured simultaneously from the science data.
- Different spectral and spatial properties of the foregrounds must be used: to exploit this, DARE will be able to gather data from 8 independent regions on the sky.
- Promising results for DARE with 3000 hours of data, but we get useful constraints with 1000 hours or less.
- The MCMC method presented here is applicable (with some modifications) to ground-based experiments, which if nothing else provide useful and stringent tests on the performance of the foreground fitting.